

Exploring the chemical enhancement for surface-enhanced Raman scattering with Au bowtie nanoantennas

David P. Fromm

Department of Chemistry, Stanford University, Stanford, California 94305

Arvind Sundaramurthy

Department of Electrical Engineering, Stanford University, Stanford, California 94305

Anika Kinkhabwala and P. James Schuck

Department of Chemistry, Stanford University, Stanford, California 94305

Gordon S. Kino

Department of Electrical Engineering, Stanford University, Stanford, California 94305

W. E. Moerner

Department of Chemistry, Stanford University, Stanford, California 94305

(Received 5 December 2005; accepted 23 December 2005; published online 10 February 2006)

Single metallic bowtie nanoantennas provide a controllable environment for surface-enhanced Raman scattering (SERS) of adsorbed molecules. Bowties have experimentally measured electromagnetic enhancements, enabling estimation of chemical enhancement for both the bulk and the few-molecule regime. Strong fluctuations of selected Raman lines imply that a small number of *p*-mercaptoaniline molecules on a single bowtie show chemical enhancement $>10^7$, much larger than previously believed, likely due to charge transfer between the Au surface and the molecule. This chemical sensitivity of SERS has significant implications for ultra-sensitive detection of single molecules. © 2006 American Institute of Physics. [DOI: [10.1063/1.2167649](https://doi.org/10.1063/1.2167649)]

Rapid and accurate detection and identification of trace amounts of chemical species is of the utmost importance in biology, chemistry, medicine, and defense. Recent advances in fluorescence spectroscopy methods offer exquisite sensitivity, enabling the ultimate in analytical detection: single molecules.¹ However, fluorescence studies require specially engineered labels, a limitation that places constraints on potential applications. Raman spectroscopy does not suffer from this limitation, since most molecules display a unique set of molecular vibrations that give rise to a distinctive chemical fingerprint, especially attractive for ultra-selective analysis.

Because Raman transitions are incredibly weak, this technique was not generally believed to offer the potential sensitivity afforded by fluorescence. However, over 30 years ago, it was first observed that the Raman signal of pyridine dramatically increases when adsorbed on a roughened Ag electrode,^{2,3} and the detailed origins of surface-enhanced Raman scattering (SERS) arising from nanostructured metals have remained a topic of debate. Researchers linked SERS signals to a combination of two effects,⁴ electromagnetic (EM) enhancement, where illumination intensity is enhanced due to sharp metal edges or plasmon effects, and chemical enhancement (CE), where the Raman cross-section of adsorbed molecules is increased above the solution value,⁵ with EM enhancement dominant. Detailed SERS experiments performed on roughened metal films in electrochemical cells revealed the importance of CE due to the applied potential dependence of SERS spectra,^{6,7} but measured values of either enhancement were not available.

Interest in the SERS mechanism blossomed with the recent observation of Raman lines apparently arising from *single molecules* adsorbed onto colloidal Ag and Au particles.^{4,8,9} To obtain the 14-order of magnitude enhancement required to make Raman signals competitive with fluorescence, it was believed that an enormous EM enhancement at “hot spots” was combined with a modest CE in these systems. Detailed investigation revealed signals are greatest at junctions between plasmonically coupled metal nanoparticles, usually formed when two or more colloidal particles of ~ 100 nm diameter happen to aggregate,^{4,10} or within fractal island films.^{11,12}

Michaels *et al.*⁴ suggested that coupled colloidal particles locally enhance the incident laser pump intensity, proportional to $|E|^2$ where E is the electric field, by as much as 10^5 . The same factor enhances the probability of detecting Raman emission as in the case of a transmitting antenna by another $|E|^2$ factor yielding a total EM enhancement potentially as high as 10^{10} .

The search for large and controllable EM enhancements has renewed interest in studying near-field coupling between metal particles, because optimized structures should readily produce SERS-active substrates. Scattering experiments have explored the resonant wavelength of subwavelength-sized particles,^{13,14} and have been accompanied by important theoretical analysis^{15,16} aiding nanostructure design. Because of their ability to increase incident and received signals of molecules in their gap, enhanced *lithographically fabricated* metallic nanostructures have been dubbed “nanoantennas” and experimentally produced large local E -fields.^{17,18}

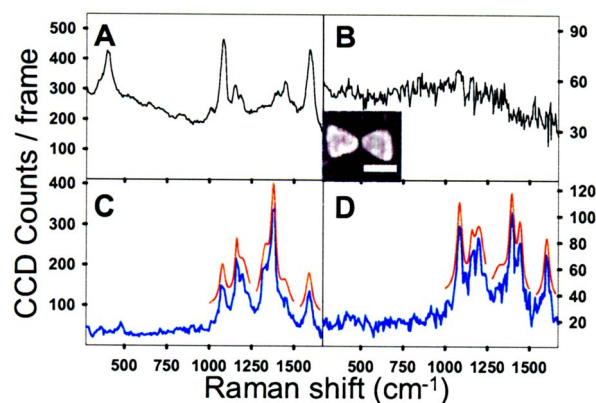


FIG. 1. (Color online) (A) Bulk SERS spectrum of *p*MA on roughened Au (10 s integration); (B) fluorescence continuum from a bowtie without adsorbed *p*MA. A representative SEM image, inset (scale bar = 100 nm, 17 nm gap); (C) SERS for *p*MA on a bowtie (2 s integration) at $t=326$ s; (D) SERS from the same bowtie at $t=441$ s. Spectra (C) and (D) differ in observed SERS modes and total intensity. Red curves in (C) and (D) are Lorentzian fits for modes between 1000 and 1600 cm^{-1} , offset for clarity. $I=155$ kW/cm^2 .

However, there is still uncertainty about the relative importance of the chemical effect, which has been estimated to enhance SERS signals by factors anywhere from 10 to as high as 10^6 .^{19,20} Here, we describe the first use of single resonant “bowtie” nanoantennas consisting of two metallic triangles oriented tip to tip designed to facilitate Raman spectroscopy, measuring SERS from chemically bonded *p*-mercaptoaniline (*p*MA) molecules. Since these structures provide a coupled plasmon system with lithographically controllable gap and known (measured) EM enhancement,¹⁷ this strategy allows the first direct exploration of the role of CE in SERS.

Au bowties were fabricated with electron-beam lithography producing gap sizes ≤ 20 nm [see Fig. 1(B) inset], and coated with *p*MA.²¹ Single bowties were pumped at their resonance (830 nm) using a linearly polarized cw laser diode, and imaged in a confocal microscope, whose collection path was equipped with a grating spectrometer and cooled Si CCD detector to acquire Raman spectra. *p*MA was chosen because the thiol moiety forms a self-assembled monolayer (SAM) upon binding to Au, it is nonabsorbing at 830 nm, and several bulk SERS studies were previously performed on *p*MA in electrochemical cells^{6,7} and on metallic nanoshells.²²

For verification of SAM preparation, a 35 nm thick roughened Au film was thermally evaporated, incubated in *p*MA,²¹ and its Raman spectrum [Fig. 1(A)] shows the peaks expected for *p*MA SERS.^{6,21} Prior to *p*MA incubation, single bowtie spectra showed only broad continuum emission [Fig. 1(B)], which scales linearly with power below the bowtie damage threshold, and is attributed to 1-photon photoluminescence excited in Au by the high optical fields near the bowtie gap.²³ Continuum is also present in the *p*MA Raman spectra on bowties and the Au film. Figures 1(C) and 1(D) show spectra (2 s integration) taken at $t=326$ s and $t=441$ s for the same bowtie revealing most of the characteristic *p*MA lines seen in Fig. 1(A). In the fingerprint region (1000–1600 cm^{-1}), 6–7 lines appear with similar Raman

shift but radically different amplitudes. For bowties, both continuum and SERS signals decrease to undetectable levels (2 s integration) when excitation light is polarized perpendicular to the bowtie long axis.

Using the known monolayer packing area of *p*MA (19 $\text{\AA}^2/\text{molecule}$)²¹ and the surface area of the bowtie experiencing enhanced *E*-fields (3000 nm^2) as predicted by electromagnetic calculations,¹⁶ the SERS signal is expected to arise from $N \sim 10^4$ molecules near the tips of the two triangles forming the bowtie gap. We estimate total SERS enhancement (TE) for *p*MA on the bowtie using: $\text{TE} = S / (N \sigma_{\text{Raman}} \Phi I)$, where S is the integrated Raman signal with continuum fluorescence subtracted ($S=5 \times 10^4$ photons/s), $\sigma_{\text{Raman}}=10^{-29}$ cm^2 , Φ the collection efficiency (0.01), and $I=6.5 \times 10^{23}$ photons/s cm^2 . Thus in Fig. 1(D), $\text{TE}=7 \times 10^7$. The TE value is often broken into EM and CE contributions in the following manner: $\text{TE} = (|E|_{\text{pump}}^2 / |E_0|^2) (|E|_{\text{transmit}}^2 / |E_0|^2) \text{CE}$, where E_0 is the incident electric field. In our bowties, $|E|_{\text{pump}}^2$ is measured¹⁷ to be $\sim 1500|E_0|^2$. Assuming that $|E|_{\text{transmit}}^2$ is also $\sim 1500|E_0|^2$, reasonable for Stokes radiation considering the fairly broad antenna resonance,¹⁴ a value of $\text{CE} \sim 30$ is obtained, consistent with previous estimates.²¹

Closer examination of the spectra shown in Fig. 1(C) and 1(D) reveals that, while the spectra are similar to the bulk spectrum [Fig. 1(A)], there are intriguing dynamics; particular lines “blink” and their mode frequencies shift discontinuously, even though an entire monolayer of *p*MA covers the bowtie. Behavior like this has been reported previously,^{4,8,9,19} but only at concentrations of a few molecules per nanoparticle. For example, in Fig. 1(C) there is an intense band at 1380 cm^{-1} that dramatically decreases in panel 1(D), where the spectrum has a lower total intensity and resembles the bulk *p*MA spectrum.

To display the typical spectral dynamics we see, Fig. 2(A) shows a “waterfall” plot, where Raman shift, $\Delta\nu$, is plotted as a function of time, and SERS intensity is displayed as a colorscale. First we consider the very intense (several thousands of photons/s) broadband flash of emission, as seen at $t \sim 150$ s in Fig. 2(A) (yellow arrow). This brief flash of broad infrared fluorescence is possibly due to the creation and subsequent destruction of a metastable Au cluster of ~ 30 atoms, previously shown to be fluorescent.^{24,25} Because these events are rare, and smaller, bluer emitting clusters would not be resonant with the bowties, we do not consider cluster formation to be central to the SERS effects reported here.

Immediately obvious is that the intensity and frequency of some Raman lines are stable, but vary significantly for other modes. Using previous *p*MA mode assignments,⁶ modes at 1077 and 1590 cm^{-1} , assigned to a_1 symmetry,²¹ are stable in their frequency, and their intensity varies by less than a factor of 2 over the entire spectral series. For this reason, we will use the a_1 modes as key reference modes in our analysis below. Conversely, peaks at 1160, 1195, 1325, 1380, and 1450 cm^{-1} , which have b_2 symmetry and have been interpreted as charge transfer (CT) modes for *p*MA,^{6,7} display large frequency and intensity fluctuations. For frequency fluctuations, Fig. 2(B) displays the center frequency

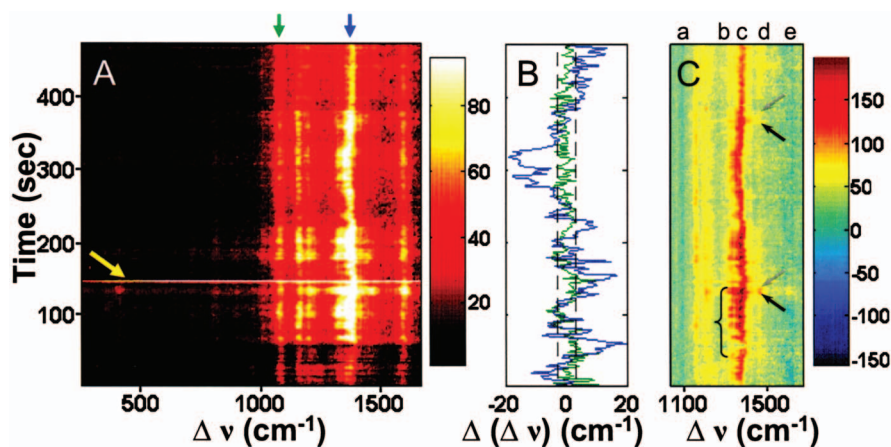


FIG. 2. (Color) (A) Waterfall plot showing the evolution of SERS spectra with time (2 s integration). $I=11 \text{ kW/cm}^2$ for $t=0-57 \text{ s}$, and $I=38 \text{ kW/cm}^2$ for $t=57-460 \text{ s}$. Notice the intense flash of continuum fluorescence at $t \sim 150 \text{ s}$ (yellow arrow), the relatively stable modes at 1077 and 1590 cm^{-1} , and fluctuations of both frequency and amplitude for the 1160 , 1195 , 1325 , 1380 , and 1450 cm^{-1} Raman modes. (B) Mode center frequency fluctuations for 1077 (green) and 1380 cm^{-1} (blue) lines, with error bar shown as dashed lines.²¹ Only modes with b_2 symmetry fluctuate significantly. (C) A waterfall plot of difference spectra showing discrete intensity fluctuations evident for b_2 modes. Note the blinking in the 1325 cm^{-1} mode (bracket, mode “b”) and “on” (black arrows) “off” (gray arrows) switching for mode “d,” 1450 cm^{-1} . Mode “c” (1380 cm^{-1}) also fluctuates. Modes (“a,” 1077 cm^{-1} and “e,” 1590 cm^{-1}) having a_1 symmetry are suppressed by the subtraction which is produced in two steps: first, intensity for the non-fluctuating fiducial 1077 cm^{-1} mode for each Raman spectrum in Fig. 2(A) is used to scale the observed bulk SERS spectrum of pMA on a rough Au film [Fig. 1(A)], and this scaled bulk SERS spectrum is subtracted from the measured spectrum in A.

time trajectory of each of two modes, determined by fitting the peaks to Lorentzian profiles.²¹ The 1077 cm^{-1} (a_1 symmetry, C–S stretch) reference mode (green) displays little change, while the 1380 cm^{-1} (b_2 symmetry, combination C–H bend and C–C stretch) mode (blue) varies significantly, often “jumping” as much as $\pm 20 \text{ cm}^{-1}$ from its expected value.

We assume that the Raman spectra include contributions from a large number of molecules with the small enhancement calculated above, plus a contribution from a small number of molecules with large (possibly huge) enhancement. To highlight the latter contribution, we use the strength of the 1077 cm^{-1} reference mode to scale the bulk SERS spectrum [Fig. 1(A)] and subtract it from the observed spectra, generating the difference spectra in Fig. 2(C). Note that scaling by the 1077 cm^{-1} mode automatically makes the other a_1 mode at 1590 cm^{-1} disappear. (It should be noted that a single blinking event is observed for mode “e” at $t \sim 110 \text{ s}$, because a weak b_2 symmetry mode exists at 1575 cm^{-1} .) Figure 2(C) shows strong on/off amplitude fluctuations in several b_2 modes (designated “b” and “d”), for example at the arrows and bracket. The discrete fluctuations of the b_2 modes suggests a dynamic CT process is occurring and that the SERS signal is quite sensitive to CT.

Due to the difficulty in synchronizing even a small number of molecules, discrete intensity fluctuations (blinking) and spectral jumping are two spectroscopic signatures of single molecules.¹ We note that blinking alone is *not necessarily* a fair indicator of single-molecule SERS, because the EM enhancement could rapidly change, leading to temporary fluctuations of SERS signals for all molecules near the gap. The stability of the reference a_1 modes shows that large EM changes are not occurring. Because the frequency and/or amplitude of the b_2 modes often discontinuously “jump,” and in any one 2 s spectrum, a single frequency is observed (movie in Ref. 21), this observation implies that there is either a

synchronization of all 10^4 molecules near the gap, very unlikely, or that a very small subset of molecules is responsible for the discontinuous changes, with the possibility that this SERS signature is from a single molecule. Therefore, the varying frequencies [Fig. 2(B)] and amplitudes [Fig. 2(C)] of the CT b_2 modes suggest that selected molecule(s) are subject to a huge variation in their CE, from ~ 30 to as much as $\sim 10^7$. Apparently, signal from a small subset is comparable, and often exceeds (by more than an order of magnitude) that of the remaining 10^4 molecules near the bowtie gap, even though both are subject to the same EM value, not an unreasonable possibility.²⁶

Because we are able to distinguish EM and CE enhancement components, and EM is not varying, it is clear that fluctuating CE is of primary importance in producing the dynamic SERS behavior, a result that should help elucidate the elusive SERS mechanism. Since this idea is quite new, it is not possible to produce a precise mechanism. However, huge local CE is only present for selected low symmetry Raman modes of the benzene ring long associated with CT processes. How can a huge change in CT select for one (or a few) molecules? We suggest a few reasonable possibilities based on previous CT theory that could produce the observed effects, in order to stimulate further experimental and theoretical exploration of CE mechanisms.

Very small Ag clusters [O(10 atoms)] have shown surprisingly large Raman signals in the absence of significant EM enhancement possibly due to extremely large molecule-cluster couplings.^{27–29} Indeed, one could imagine that an Au cluster could form in the bowtie gap, and due to its small surface area, accommodate a few molecules that would be strongly enhanced, but we rarely observe cluster fluorescence or huge a_1 mode changes and thus rule out clusters as the mechanism. Helpfully, CT effects have been characterized in pMA SERS experiments on rough Ag/Au films in electrochemical cells for over a decade,^{6,7} where b_2 modes were

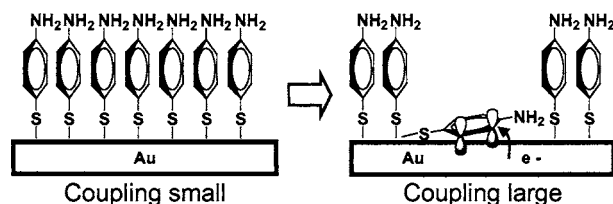


FIG. 3. Proposed CT mechanism. At a defect site, a *p*MA molecule can adsorb flat on the surface, enabling π -electrons in the benzene ring to more strongly interact with the Au surface, increasing vibronic coupling for selected modes, increasing SERS.

found to be very sensitive to applied potential, and attributed to CT. To explain this, Osawa *et al.*⁶ borrow from the comprehensive theory of Lombardi *et al.*³⁰ who describe the Herzberg-Teller effect for CT. In short, for a molecule adsorbed onto a metal, vibronic mixing between the electronic and vibrational wave functions of adsorbed molecules and those of the metal leads to electron transfer, strongly increasing Raman scattering associated with mixed modes by inducing an effect analogous to resonant Raman.²¹ A simple physical picture is presented in Fig. 3 that could produce these effects. Normally, *p*MA molecules should orient perpendicular to the surface in a close-packed SAM, making it difficult for a molecule to access the metal surface. However, at a defect, which could be light-induced, the SAM could be interrupted, allowing a single or a few *p*MA molecules to lie flat on the surface and undergo CT. In this position, benzene ring π -orbitals better couple to the Au surface, dramatically increasing vibronic mixing for selected benzene modes while leaving mixing for other modes, such as the C–S stretch, unchanged, strongly enhancing CE for selected modes only.²⁶ The transient “on”/“off” switching shown for b_2 modes in Fig. 2(C) (especially “*b*” and “*d*” lines) is consistent with this picture.

In summary, SERS is observed for *p*MA adsorbed on lithographically produced single bowtie nanoantennas with known local field enhancement, yielding the expected Raman lines on top of a fluorescent continuum originating in the Au. Despite the fact that $\sim 10^4$ molecules are present near the bowtie gap, the Raman spectra display discrete blinking and spectral shifting events, indicating that a few or possibly one molecule produces a SERS signal comparable to that of the remaining molecules. This means some molecules experience a larger CE than previously believed possible, as large as 10^7 . This is possibly understood using existing CT theory, in which we suggest that an increase in wave function mixing between π -orbitals in the benzene ring and Au occurs when a *p*MA molecule lies down at the surface. Due to the extreme sensitivity of SERS enhancement to the properties of adsorbed chemical species, the design and production of nanoantenna-based SERS substrates holds great promise in

performing trace chemical detection, but achieving the proper orientation for large CE will be a challenge.

This work was supported in part by the National Institutes of Health Grant No. 1R21-GM065331, and the Stanford Nanofabrication Facility (a member of the National Nanotechnology Infrastructure Network), supported by the National Science Foundation Grant No. ECS-9731293.

- ¹W. E. Moerner and M. Orrit, *Science* **283**, 1670 (1999).
- ²M. Fleischmann, P. J. Hendra, and A. J. McQuillan, *Chem. Phys. Lett.* **26**, 163 (1974).
- ³D. L. Jeanmaire and R. P. Van Duyne, *J. Electroanal. Chem.* **84**, 1 (1977).
- ⁴A. M. Michaels, J. Jiang, and L. Brus, *J. Phys. Chem. B* **104**, 11965 (2000).
- ⁵M. Moskovits, *Rev. Mod. Phys.* **57**, 783 (1985).
- ⁶M. Osawa, N. Matsuda, K. Yoshii, and I. Uchida, *J. Phys. Chem.* **98**, 12702 (1994).
- ⁷W. Hill and B. Wehling, *J. Phys. Chem.* **97**, 9451 (1993).
- ⁸K. Kneipp, Y. Wang, H. Kneipp, L. T. Perelman, I. Itzkan, R. R. Dasari, and M. S. Feld, *Phys. Rev. Lett.* **78**, 1667 (1997).
- ⁹S. Nie and S. R. Emory, *Science* **275**, 1102 (1997).
- ¹⁰H. Xu and M. Kaell, *ChemPhysChem* **4**, 1001 (2003).
- ¹¹S. Gresillon, L. Aigouy, A. C. Boccarda, J. C. Rivoal, X. Quelin, C. Desmarest, P. Gadenne, V. A. Shubin, A. K. Sarychev, and V. M. Shalaev, *Phys. Rev. Lett.* **82**, 4520 (1999).
- ¹²M. I. Stockman, L. N. Pandey, L. S. Muratov, and T. F. George, *Phys. Rev. Lett.* **72**, 2486 (1994).
- ¹³C. L. Haynes, A. D. McFarland, L. Zhao, R. P. Van Duyne, G. C. Schatz, L. Gunnarson, J. Prikulis, B. Kasemo, and M. Kaell, *J. Phys. Chem. B* **107**, 7337 (2003).
- ¹⁴D. P. Fromm, A. Sundaramurthy, P. J. Schuck, G. S. Kino, and W. E. Moerner, *Nano Lett.* **4**, 957 (2004).
- ¹⁵E. Hao and G. C. Schatz, *J. Chem. Phys.* **120**, 357 (2004).
- ¹⁶A. Sundaramurthy, D. P. Fromm, K. B. Crozier, P. J. Schuck, W. E. Moerner, and G. S. Kino, *Phys. Rev. B* **72**, 165409 (2005).
- ¹⁷P. J. Schuck, D. P. Fromm, A. Sundaramurthy, G. S. Kino, and W. E. Moerner, *Phys. Rev. Lett.* **94**, 017402 (2005).
- ¹⁸P. Muhlschlegel, H. J. Eisler, O. J. F. Martin, B. Hecht, and D. W. Pohl, *Science* **308**, 1607 (2005).
- ¹⁹G. Haran, *Isr. J. Chem.* **44**, 385 (2004).
- ²⁰A. Otto, *J. Raman Spectrosc.* **33**, 593 (2002).
- ²¹See EPAPS Document No. E-JCPSA6-124-009604 for a movie of time-dependent SERS spectra, details regarding experimental conditions and statistical methods employed, and for detailed discussion. This document can be reached via a direct link in the online article’s HTML reference section or via the EPAPS homepage (<http://www.aip.org/pubservs/epaps.html>).
- ²²J. B. Jackson and N. J. Halas, *Proc. Nat. Acad. Sci. USA* **101**, 17930 (2004).
- ²³M. R. Beversluis, A. Bouhelier, and L. Novotny, *Phys. Rev. B* **68**, 115433 (2003).
- ²⁴S. Link, A. Beeby, S. FitzGerald, M. A. El-Sayed, T. G. Schaaff, and R. L. Whetten, *J. Phys. Chem. B* **106**, 3410 (2002).
- ²⁵J. Zheng, C. Zhang, and R. M. Dickson, *Phys. Rev. Lett.* **93**, 077402 (2004).
- ²⁶Z. Wang and L. J. Rothberg, *J. Phys. Chem. B* **109**, 3387 (2005).
- ²⁷M. L. Jacobson and K. L. Rowlen, *Chem. Phys. Lett.* **401**, 52 (2005).
- ²⁸O. L. A. Monti, J. T. Fourkas, and D. J. Nesbitt, *J. Phys. Chem. B* **108**, 1604 (2004).
- ²⁹L. P. Capadona, J. Zheng, J. I. Gonzalez, T. H. Lee, S. A. Patel, and R. M. Dickson, *Phys. Rev. Lett.* **94**, 058301 (2005).
- ³⁰J. R. Lombardi, R. Birke, L. T. Lu, and J. Xu, *J. Chem. Phys.* **84**, 4174 (1986).